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Citation: *Journal of Applied Physics* **100**, 113904 (2006); doi: 10.1063/1.2390622

View online: <http://dx.doi.org/10.1063/1.2390622>

View Table of Contents: <http://aip.scitation.org/toc/jap/100/11>

Published by the *American Institute of Physics*

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# Magnesium diboride superconductor thin film tunnel junctions for superconductive electronics

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(Received 24 August 2006; accepted 19 September 2006; published online 7 December 2006)

Based on superconducting MgB<sub>2</sub> films with higher critical temperature of 39 K and the advantage of the conventional superconductors, those that follow Bardeen-Cooper-Shrieffer theory, fabrication of quasiparticle, and Josephson tunnel junctions have been investigated. To explore the potential of MgB<sub>2</sub> for superconductive electronics, the essentials such as clean high quality thin film surfaces and reproducible tunnel junction fabrication are addressed. Our results show clean tunneling characteristics for *in situ* prepared MgB<sub>2</sub> junctions on Si wafer with the measured superconducting energy gap values in good agreement with theory and the feasibility of the technology. The recent results on all-epitaxial MgB<sub>2</sub>/MgO/MgB<sub>2</sub> Josephson junctions will be also reported. © 2006 American Institute of Physics. [DOI: 10.1063/1.2390622]

## INTRODUCTION

After the discovery of superconductivity in the intermetallic compound magnesium diboride (MgB<sub>2</sub>) with a high critical temperature of 39 K,<sup>1</sup> a number of groups have investigated MgB<sub>2</sub> due to its great potential for superconductive magnets and electronic applications. In comparison with the conventional metallic superconductors, MgB<sub>2</sub> has the potential for higher operating speed due to its larger energy gap. For the realization of superconducting electronics (SCE) the availability of high quality thin films and the technology to fabricate Josephson circuits based on these films are essential.<sup>2,3</sup>

The growth of MgB<sub>2</sub> thin films, however, was challenging and a number of problems had to be solved, especially in order to obtain smooth and single-phase films suitable for multilayer structures that are important for the application in superconducting electronics. The problems such as high volatility of magnesium, oxygen impurities, and surface roughness had to be overcome.<sup>4</sup> Due to the high vapor pressure of Mg, the phase stability of MgB<sub>2</sub> with high crystallinity is the most serious problem. Therefore, earlier two-step growth, in other words, *ex situ* postannealing was developed. Later, one-step growth, namely, *in situ* physical vapor deposition, such as molecular beam epitaxy (MBE) was reported.<sup>2,4</sup> The detailed procedure for the growth of single-phase epitaxial MgB<sub>2</sub> films on Si(111) substrate has been described in our earlier papers.<sup>4,5</sup>

The electronic structure and superconducting (SC) behavior of MgB<sub>2</sub> are theoretically well described using the conventional electron-phonon mechanism.<sup>6,7</sup> Anisotropic SC energy gaps,  $\Delta_\sigma$  for the two quasi-two-dimensional  $\sigma$  bands and  $\Delta_\pi$  for the pair of three-dimensional  $\pi$  bands, ascribed to two different parts of the Fermi surface, were predicted

theoretically.<sup>8</sup> The calculation by Choi *et al.*<sup>9</sup> showed that the electronic states are dominated by orbitals in the boron plane coupling strongly to specific phonon modes giving rise to multiple gaps. Those authors predicted SC energy gaps from 6.4 to 7.2 meV on the  $\sigma$  band and from 1.2 to 3.7 meV on the  $\pi$  band. Although the multiband picture of superconductivity in MgB<sub>2</sub> has been studied to determine the SC energy gaps by various surface sensitive techniques such as probe tunneling, point contact Andreev reflection, high resolution angle resolved photoelectron spectroscopy, etc.,<sup>10-16</sup> leakage free planar tunneling junctions have provided the most reliable SC gap information in conventional superconductors.<sup>5,17</sup> Indeed, not only for fundamental studies but also for many superconductor electronic applications of MgB<sub>2</sub>-based junctions, it is important to know how the two-band superconductivity will play a role. Besides, the pertinence of MgB<sub>2</sub> for SCE strongly depends on the possibility of making well-controlled Josephson junctions. In this paper our recent results on the SC characteristics of epitaxial MgB<sub>2</sub> films using MgB<sub>2</sub>-based tunnel junctions and the fabrication of Josephson junction, in particular, all-MgB<sub>2</sub> junctions (such as MgB<sub>2</sub>/MgO/MgB<sub>2</sub>) are reported.

## SAMPLE PREPARATION

MgB<sub>2</sub> thin film planar junctions were prepared, all *in situ*, using shadow masks in a MBE system with the base pressure in the 10<sup>-10</sup> Torr range. Onto an etched Si(111) substrate, 50 Å of MgO was grown by e-beam evaporation at a substrate temperature ( $T_s$ ) of 563 K. This was followed by the coevaporation of Mg (from a Knudsen cell) and B (from an e-beam source) with independent rate controllers to obtain MgB<sub>2</sub> films with thickness of 560 Å. The optimum conditions required to obtain high quality MgB<sub>2</sub> films<sup>4</sup> and to fabricate MgB<sub>2</sub>-based tunnel junctions with Al<sub>2</sub>O<sub>3</sub> tunnel barriers (MgB<sub>2</sub>/Al<sub>2</sub>O<sub>3</sub>/V) (Ref. 5) are described in our earlier studies. MgB<sub>2</sub>/MgO/MgB<sub>2</sub> junctions were fabricated

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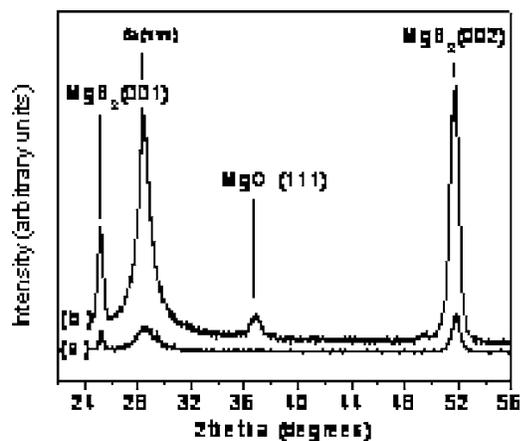


FIG. 1. XRD  $\theta$ - $2\theta$  patterns of MgB<sub>2</sub> thin film grown at 563 K (a), and of MgB<sub>2</sub>/MgO/MgB<sub>2</sub> trilayer with 100 Å thick MgO film (b). 50 Å of MgO was deposited beyond these films as a capping layer to avoid the oxygen contamination.

slightly differently; MgO sandwiched between the two MgB<sub>2</sub> layers is the tunnel barrier. To cover the edges of the bottom MgB<sub>2</sub> film strips, a 70-Å-thick MgO film was deposited, thus defining the tunneling width (80  $\mu$ m) of the bottom electrode. Cross strips of 560-Å-thick MgB<sub>2</sub> were again deposited as the top electrode, keeping the  $T_S$  of 563 K throughout. These junctions were protected with 50 Å of MgO film. The thickness of the barrier was varied to obtain tunnel junctions with an optimum resistance ( $R_J$ ) in the range of a few ohms to a few kilohms.

## RESULTS

The structural characterization of our MgB<sub>2</sub> films by x-ray diffraction show epitaxial growth, see Fig. 1(a). Spectrum (b) in Fig. 1 shows that we succeeded in obtaining as-grown high-quality MgB<sub>2</sub>/MgO/MgB<sub>2</sub> epitaxial superconducting trilayer structure at our growth conditions over Si(111) with a 50 Å thick MgO seed layer. It is notable that the top MgB<sub>2</sub> film could be grown epitaxially on a MgO(111) barrier layer grown over a MgB<sub>2</sub> bottom film. One can see clearly that the peak intensity of MgO(111) becomes stronger as the thickness of MgO increases from 50 to 100 Å. It may be noted that the top MgB<sub>2</sub> film can grow with good epitaxy directly onto the MgO(111) layer because the lattice mismatch is only a  $\sim 2\%$  for two unit cells of MgB<sub>2</sub> on a MgO(111) unit cell. Critical temperatures ( $T_c$ ) of 28 and 33 K were obtained for the top and the bottom MgB<sub>2</sub> films, respectively. The lower  $T_c$  of the top MgB<sub>2</sub> film may be attributed to the poor microstructure at the MgO(111)/MgB<sub>2</sub> interface. The  $T_c$  needs to be further optimized.

The tunneling conductance for a MgB<sub>2</sub>/Al<sub>2</sub>O<sub>3</sub>/V junction taken at 0.45 K is shown in Fig. 2. The zero conductance region is very well established up to nearly  $\pm 2.5$  mV, and the main peaks are noticeably sharper, as expected for a good quality superconductor/insulator/superconductor (SIS) tunnel junction. The SC gaps ( $2\Delta$ ) of both the MgB<sub>2</sub> and the V films are determined unambiguously from the sum of conductance peaks. From the simple analysis, the  $\Delta$  value for V

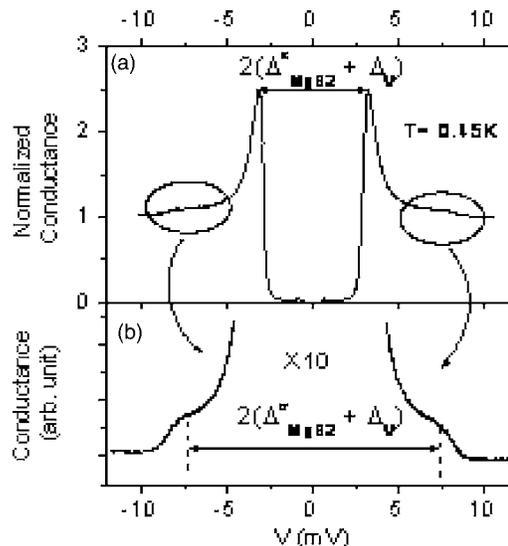


FIG. 2. Tunneling characteristics at lower temperature (0.45 K) showing not only the SC gap energy sum peaks due to  $\pi$  bands in MgB<sub>2</sub> and V but also those due to the  $\sigma$  band of MgB<sub>2</sub> and V.

is 0.70 meV, whereas for MgB<sub>2</sub>, the  $\Delta$  value is 2.4 meV, which we attribute to the  $\pi$  band ( $\Delta_{\pi}$ ).<sup>8,9</sup> Careful observation of this plot reveals additional features at voltages slightly higher than the main peaks as well as tiny voltage steps on either side of  $V=0$ . In the bottom part of this figure [see Fig. 2(b)], at higher voltages the conductance is amplified ten times in order to show the less well resolved second peaks (rather small in amplitude compared to the main peaks). The position of the second set of peaks is determined to be  $\pm 7.5$  meV by taking the derivative of this curve, from which we obtain the  $\sigma$  gap of MgB<sub>2</sub>, as  $\Delta_{\sigma}=6.8$  meV, which is due to the  $\sigma$  band (after subtracting the V SC gap). The tiny voltage steps (at  $\pm 0.6$  mV) are attributed to the presence of a small amount of Josephson pair current in the junction. When we make the barrier thinner, we clearly see a large Josephson pair current, as shown in Fig. 3 for one such device. A current hysteresis, as well as self-induced voltage steps, was also observed in such structures with thinner barriers.

We have also seen Josephson pair tunneling with MgO barriers, similar to that seen with Al<sub>2</sub>O<sub>3</sub> barriers. MgO barriers were deposited from an e-gun source with a  $T_S$  at 563 K. The inset of Fig. 4 shows that we successfully grew all-epitaxial MgB<sub>2</sub>/MgO/MgB<sub>2</sub> junctions with excellent SC properties. The critical current was seen to be close to the

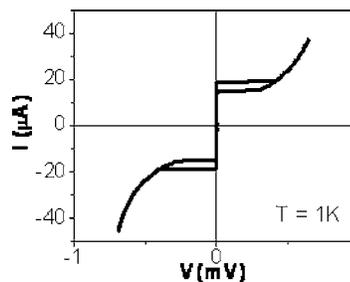


FIG. 3. Josephson pair current observed at zero bias for a MgB<sub>2</sub>/Al<sub>2</sub>O<sub>3</sub>/V junction with a thinner Al<sub>2</sub>O<sub>3</sub> tunnel barrier.

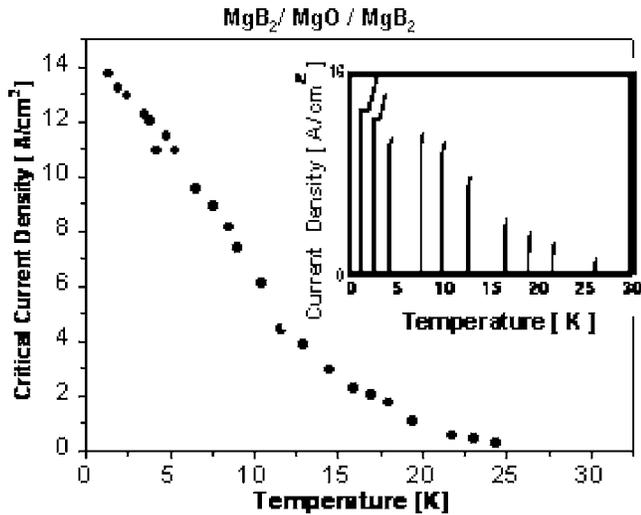


FIG. 4. Temperature dependence of Josephson current density in  $\text{MgB}_2/\text{MgO}/\text{MgB}_2$  junction with  $\sim 200 \mu\text{m}^2$  junction area; Josephson current as a function of temperature for all  $\text{MgB}_2$  tunnel junctions ( $\text{MgB}_2/\text{MgO}/\text{MgB}_2$ ), as shown in the inset.

value of  $T_c$  of the top  $\text{MgB}_2$  layer (28 K). A similar behavior was also observed in  $\text{MgB}_2/\text{Al}_2\text{O}_3/\text{MgB}_2$  junctions with 20-Å-thick  $\text{Al}_2\text{O}_3$  barriers. Figure 4 shows that the temperature-dependent critical current density does not follow the Ambegaokar-Baratoff theory.<sup>18</sup> This behavior in our junctions could result from the excess of Mg atoms between  $\text{MgB}_2/\text{MgO}$  interfaces. One of the important criteria for Josephson junction application is the product  $I_C R_N$  where  $I_C$  is the Josephson critical current and  $R_N$  is the normal state resistance of the junction above the gap voltage. For a given operating frequency, it is beneficial to have as high an  $I_C R_N$  value as possible.<sup>19</sup> An  $I_C R_N$  value of  $\sim 4$  meV was obtained for our Josephson junctions, which is comparable with  $\pi$ -band gap energy. This parameter is proportional to the SC gap energy, and as such, having two large gap values,  $\Delta_\pi$  and  $\Delta_\sigma$ , even at temperatures above 20 K, has a distinct advantage for high-speed electronics applications compared to conventional low  $T_c$  superconductors.

In summary, we investigated the SC characteristics of epitaxial  $\text{MgB}_2(001)$  films using quasiparticle tunneling with

a clear determination of two SC energy gaps for the  $\pi$  and the  $\sigma$  bands. Furthermore, the feasibility of the fabrication in a well-controlled manner of all  $\text{MgB}_2$  Josephson junctions, as well as tunnel junctions with two different barriers was reported with excellent junction quality consistent with theoretical predictions. Our work on all  $\text{MgB}_2$  junctions on Si wafers might open up possibilities for  $\text{MgB}_2$ -based high temperature superconductive electronics.

## ACKNOWLEDGMENTS

This research was supported by an ONR Grant no. N00014-02-1-0119 and partly by the Basic Research Program of the Korea Science and Engineering Foundation (No. R01-2006-000-11227-0) and a Korea Research Foundation grant (KRF-2006-531-C00026) funded by the Korean Government.

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